



Improvement of Gamma Radiation Resistance in CeO_2 Doped Phosphate Glass by Co-doping with Sb_2O_3

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Abstract: It has been established that the gamma radiation resistances of innovative type of phosphate glass can be greatly enhanced by CeO_2 and Sb_2O_3 co-doping. With the doping of CeO_2 , the radiation resistance (transmittance decrease ratio) is enhanced from 57.39% to 73.9% at 525 nm, and from 56.4% to 61.9% at 385 nm, respectively, when optical glasses were exposed to the gamma radiation with the dose of 250 krad (Si). It further increases to 92.4% at 525 nm by co-doping with Sb_2O_3 , meanwhile, the induced optical losses were distinctly restrained at 1064 nm and 1550 nm, which shows potential applications in the fields of space-born star camera systems, laser window optics, fiber gyroscopes and communications.

Key words: Dope, Gamma Radiation, Resistance

1. Introduction

Radiation damage, especially induced optical losses in glass [1], are being actively studied by numerous researchers for carrying out the mechanism of optical losses in the past few years [2–5], to expand the corresponding applications in a wide range of fields including space exploring and communications [6–8]. It is well known that the induced optical loss is mainly associated with the formation of color centers in glass matrix [3, 9, 10]. The radiation resistance [11,12], i.e. the suppression of the increase in optical losses, of the optical glasses, protective filters, color-separation gratings and ytterbium-doped laser fibers can be improved with the doping of cerium [13,14]. The cerium ions are found in these optical glasses with the polyvalent ion of Ce^{3+} and Ce^{4+} , Ce^{3+} ions are converted to Ce^{4+} by capturing the radiation-induced holes and then inhibit the formation of centers associated with trapping holes. And Ce^{4+} inhibit the formation of centers resulting from trapped electrons, which prevents the formation of color centers that absorb in the visible range [12, 15].

Silicate glasses are usually chosen as host materials to incorporate Ce ions to improve their radiation resistance for use in space-born optical systems to guarantee that the required system performance is maintained during long-term mission times [16, 17]. For example, the optical systems of long-life star cameras/trackers for satellites demand the anti-radiation optical materials with higher radiation stability in the UV to visible spectral range, i.e. smaller transmittance loss due

to space radiations, in order to obtain high quality star images [18]. We, previously, have reported the high laser induced damage threshold in fluoride -containing phosphate based glasses [19], compared with the silica glasses, which can be used in the high energy laser facilities, such as US NIF and SG III in China. Combining the advantage of fluoride-containing phosphate glass, a group of CeO_2 doped phosphate glasses have been investigated, possessing a high radiation resistance with a broad transmission window of 0.4-1.6 μm , and easily fiberized, i.e., a new group of optical window glasses in the near IR range at around 1.0 μm and 1.5 μm .

In this work, we show that the addition of antimony (Sb_2O_3) to the glass composition can further improve the radiation resistance in Ce ions doped optical glass, and maintain the high transmittance at 1064 nm, 1550 nm, and smaller transmittance loss in the short wavelength range (400-780 nm) after being exposed to the gamma radiation, which shows the potential applications in high energy lasers, space laser communications and space-born optical systems, respectively.

2. Experimental

A group of un-doped, cerium (0.11 wt%) doped, cerium (0.11 wt%) + silicon (0.1 wt%) co-doped, and cerium (0.11 wt%) + antimony (1.28 wt%) co-doped samples made from the phosphate host glass with a weight composition (wt%) of Li_2O (0.5-2), K_2O (3-5), MgO (3-5), BaO (7-10), Al_2O_3 (8-11), P_2O_5 (59-64), SiO_2 (0-1.5) and H_3BO_3 (0-1.5) have been prepared and studied, and labeled as G, G: Ce, G: Ce, Si and G: Ce, Sb, respectively. The raw materials ($\geq 99.99\%$, Fe, Ni, Co, Cu < 1 ppm) were weighted and mixed in a silica crucible and then melted in an electric furnace through a mechanical stirring process at 1200 $^\circ\text{C}$ for 2 h under ambient atmosphere and then cast into a copper mold preheated at 300 $^\circ\text{C}$. The molded samples were annealed at 400 $^\circ\text{C}$ (near the glass transition temperature) through a precision annealing process with a cooling rate of -1 $^\circ\text{C}/\text{h}$ from 400 $^\circ\text{C}$ to 300 $^\circ\text{C}$, followed by -3 $^\circ\text{C}/\text{h}$ from 300 $^\circ\text{C}$ to 150 $^\circ\text{C}$ [19]. After annealing, all the glass samples were cut and precisely polished using zirconia (ZrO_2) micro-particles into the size of 20 mm \times 20 mm \times 2 mm. Besides, a series of experiments were done to characterize the produced glasses.

All of the glass samples were exposed to gamma radiation using a ^{60}Co source at a radiation dose rate of 41.79 rad (Si)/s to accumulate absorbed doses of 100 krad (Si) and 250 krad (Si), respectively. The transmission spectra were recorded with a UV-VIS-NIR spectrophotometer (Shimadzu UV-3101) in the range of 200-1600 nm for all the samples with a thickness of 2 mm. Raman spectra were collected with a Jobin-Yvonne LabRam microscope with a 532 nm laser excitation in the range of 100-1500 cm^{-1} . The X-ray photoelectron spectroscopy measurements were conducted on a Thermo Advantage X-ray photoelectron spectrometer (XPS) at room temperature using Al K (1486.6 eV) as the radiation source. XPS spectra were collected on a vacuum cleaved sample, and the vacuum pressure was 5×10^{-9} torr. The surfaces charging was corrected using the C1s level at 284.8 eV, as an internal standard. All the samples were sputtered for 30s to remove carbon pollution. EPR spectra were measured at 100 K using a Bruker A300 spectrometer operating in the X-band frequency (9.44 GHz) with field modulation 100 kHz. The microwave power used was 0.2 mW.

3. Results and discussion

The optical transmission spectra of un-doped and doped phosphate based glasses are compared in Fig. 1(a). It can be found that all of the samples keep a transmittance higher than 91.7% in the wavelength range of 350-1350 nm, even cerium were doped. Silicon was introduced to produce a reducing glass melting atmosphere that help decrease the POHC and Fe^{3+} related defects in the

phosphate based glass [19], which directly resulted in the blue-shift of the ultraviolet absorption edge of the glass. It is observed that the UV absorption edge of these Ce doped samples show an obvious red-shift as compared with the Ce-free sample, which is associated with the absorption of Ce^{4+} ions in glass matrix. And the corresponding absorption spectra of these glasses are illustrated in Fig. 1(b). The strongest peaks locating at around 4.3 eV arising from the 4f-5d transitions of Ce^{3+} ions [20]. The bands cover a broad region (4.8-6.2 eV) which is demonstrated by several types of intrinsic defects in phosphate glass, especially for the Ce-free sample (G). As for the cerium-containing samples, the bands in the high energy region show a rather rapid increase, which is associated with the absorption of Ce^{4+} ions. It can be noticed that the introduction of Si shows significant reduction effect. Which result in the slight increase of the peak at around 4.3 eV, and obvious decrease of the bands at 4.8-6.2 eV.

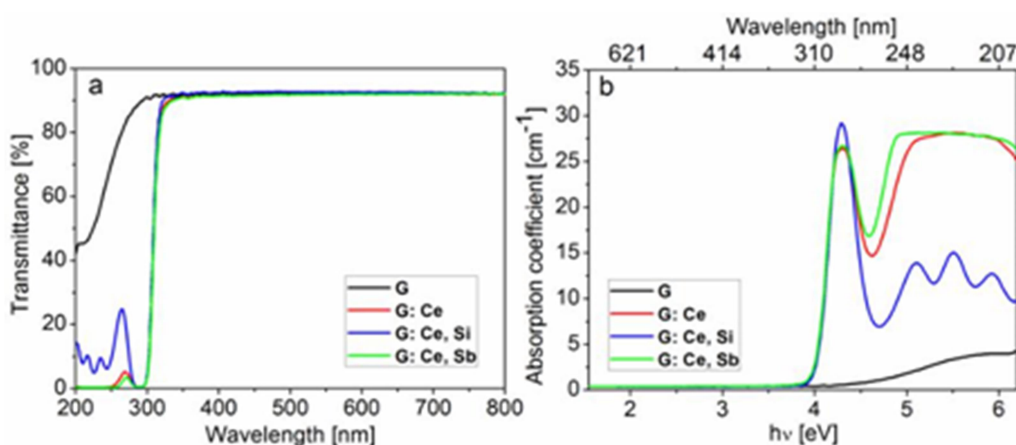


Fig. 1. (a) Measured transmission spectra, (b) calculated absorption spectra of G, G: Ce, G: Ce, Si and G: Ce, Sb glasses.

Figure 2 compares the transmission spectra of these glasses before and after exposure to the gamma radiation of 100 krad (Si) and 250 krad (Si), respectively. Gamma radiation leads to an obvious decline of the transmittance in the UV and visible spectrum range, indicating that more color centers were generated in these glasses, as shown the inset (Fig. 2(a)). With the increase of the radiation dose, the transmittance further decrease, especially for the absorption bands at around 385 nm and 525 nm, which is associated with the phosphate-related non-bridging oxygen hole center (POHC) and electron centers trapped on the central phosphorus atom in PO_3 group (PO_3 -EC) defects [21,22]. As for the cerium single- doped glass, to some extent, the radiation resistance (transmittance ratio of irradiated sample to the non-irradiated sample at certain wavelength) is improved from 57.39% to 73.9% at 525 nm, and from 56.4% to 61.9% at 385 nm, when it was exposed to gamma radiation with the dose of 250 krad (Si). And the above-mentioned absorption was also restrained in the co-doped samples, especially for the co-doped G: Ce, Sb sample (the radiation resistance increases to 92.4% at 525 nm by co-doping with Sb_2O_3). This is in accordance with the changes of glass's color in Fig. 2 (inset). In order remove the effect of indeed surface finish, we also present the absorbance after deducting the intrinsic absorbance of G, G: Ce, G: Ce, Si and G: Ce, Sb as shown in Fig. 3. It can be found that the absorption decreases at 2.36 eV (525 nm) and 3.22 eV (385 nm) by co-doping with Sb_2O_3 in Ce doped sample. These results indicate that the gamma radiation resistance can be significant improved by co-doping with Sb^{3+} .

Herein, as we further increase cerium concentration to 0.66 wt % (0.75 mol %) (Antimony is kept the same 1.28 wt% (0.5 mol %)) in the co-doped sample (G: Ce, Sb-2), its radiation resistance further increases to 99.3% at 525 nm and 82.5% at 385 nm, as shown in Fig. 4(a). And the absorption spectra after deducting the intrinsic spectra of G: Ce, Sb-2 is also presents in Fig. 4(b). The absorption at around 2.36 eV (525 nm) and 3.22 eV (385 nm) are 0.94 cm^{-1} and 0.034 cm^{-1} , respectively. Which is better than that results for the sample with same Sb_2O_3 doping concentration (0.5 mol %) reported made by Xiao et al. [23].

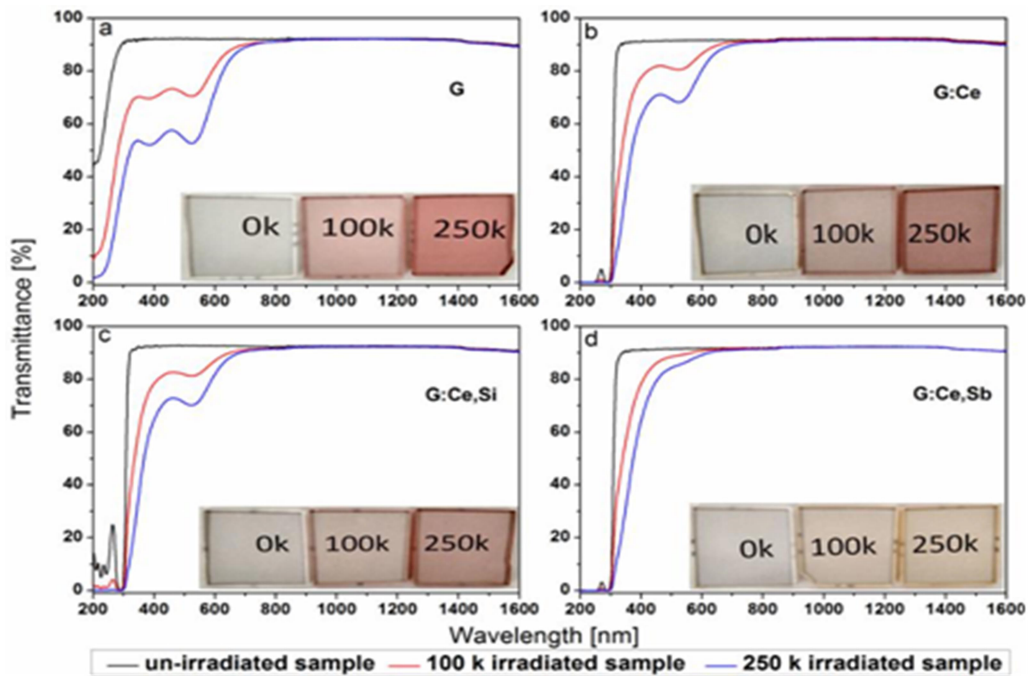


Fig. 2. Measured transmission spectra and (Inset) photograph of (a) G, (b) G: Ce, (c) G: Ce, Si and (d) G: Ce, Sb glasses before and after gamma radiation (100 and 250 krad (Si), respectively).

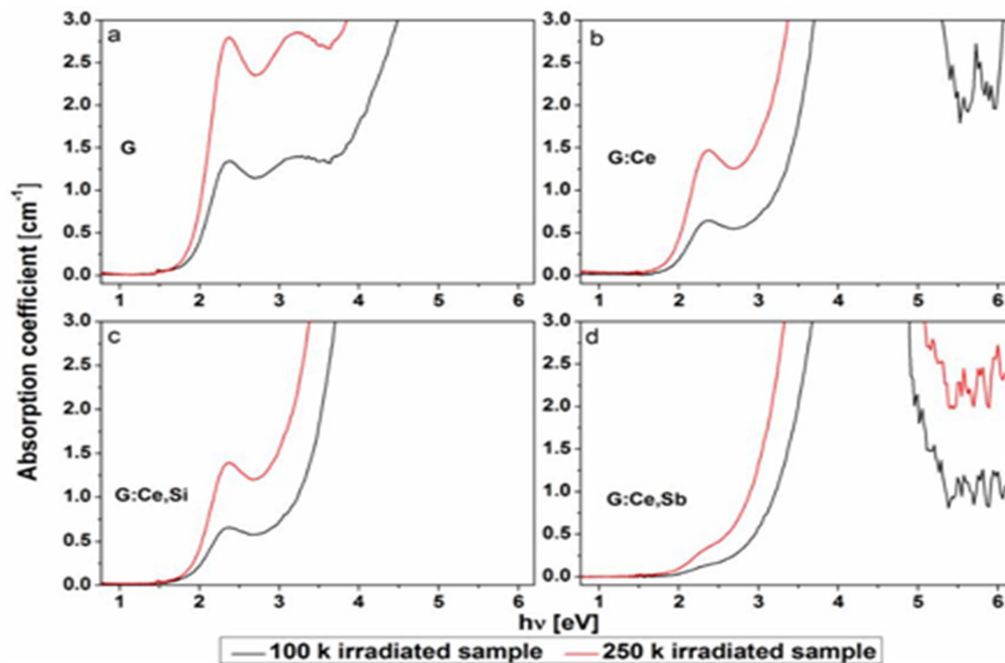


Fig. 3. The absorption spectra after minus the intrinsic absorbance of (a) G, (b) G: Ce, (c) G: Ce, Si and (d) G: Ce, Sb glasses after gamma radiation (100 and 250 krad (Si), respectively).

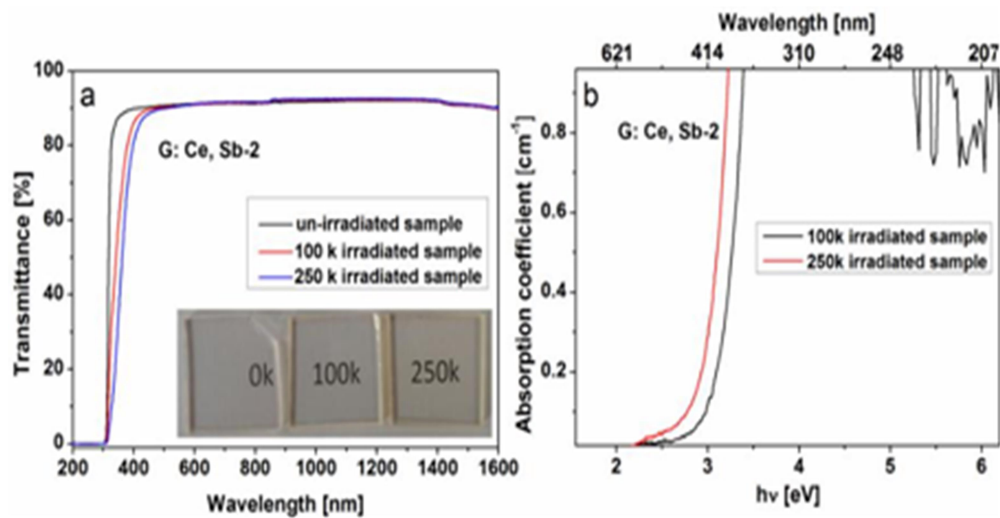


Fig. 4. (a) Measured transmission spectra and (Inset) photographs of G: Ce, Sb-2 (higher Ce doping concentration) glasses before and after gamma radiation (100, 250 krad (Si), respectively), and (b) the absorption spectra after deducting the intrinsic absorbance of G: Ce,Sb-2 glass after gamma radiation (100 and 250 krad (Si), respectively).

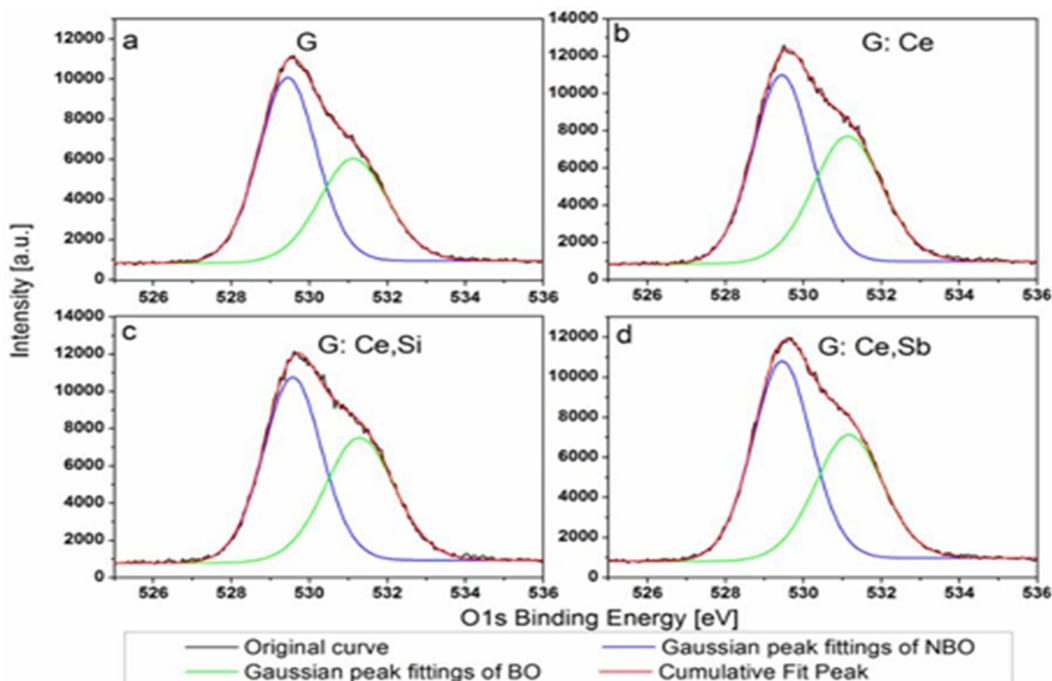


Fig. 5. Measured O1s XPS spectra of (a) G, (b) G: Ce, (c) G: Ce,Si and (d) G: Ce,Sb glasses with Gaussian peak fittings (blue and green curves represent NBO and BO, respectively).

The O1s XPS spectra are best fitted with two Voigt peaks and given in Fig. 5. These bands peaking at the lower and higher binding energy are assigned to non-bridging oxygen [24] (NBO) bonding to the glass modifier ions, and bridging oxygen (BO) bonding to the glass former ions, respectively. As one can see that the Ce doped sample (G: Ce) shows an obvious decrease of NBO (the ratios of the peak's area of NBO to the sum peak's area decrease from 58.7% to 54.6%), and an increase of BO, compared with the Ce-free sample (G). It indicates that the doping of Ce contributes to the breakage of NBO bond. However, as for the Ce, Sb co-doped sample (G: Ce, Sb), the band's area of NBO increase from 55.1% to 57.6% and BO's band area decrease compared with Ce, Si co-doped sample (G: Ce, Si). These results suggest that the broken bond of NBO bond can was fixed with co-doping of Si and Sb in Ce-containing phosphate glasses.

Figure 6 shows the typical Raman spectra of the series of glass samples as we mentioned before [25]. The vertical coordinates of spectra including G: Ce, G: Ce, Si and G: Ce, Sb samples were increased by 6400, 12800 and 19200 a.u., respectively. The obvious variations at around 1262 cm^{-1} are assigned to asymmetric stretching of non-bridging oxygens [26]. It can be found that the band intensity of 1262 cm^{-1} , as doping of Ce in glasses, decrease and the band increase in these co-doped samples, especially for the co-doped G: Ce, Sb sample. This is in accordance with the XPS spectra.

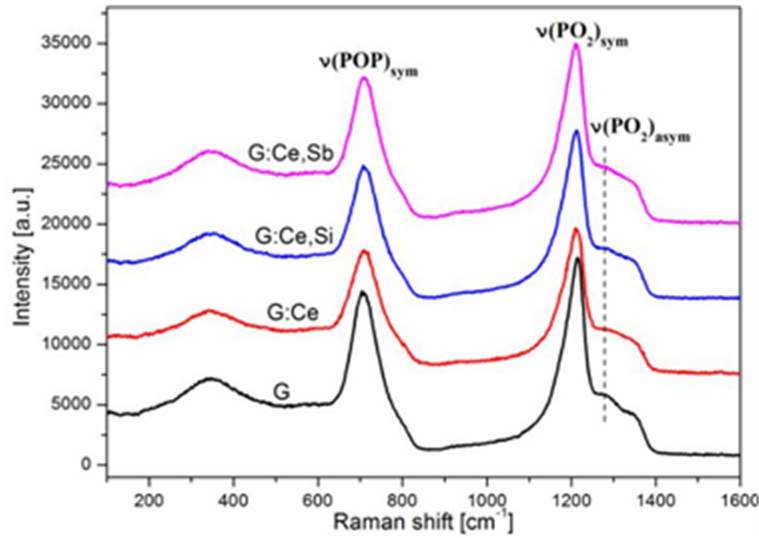


Fig. 6. Measured Raman spectra of G, G: Ce, G: Ce, Si and G: Ce, Sb glasses.

Typical EPR spectra of the above mentioned samples before and after gamma radiation are compared in Fig. 7. The EPR spectra of non-irradiated samples dominated by the signals with the magnetic field around 3355 G and 3240-3500 G, which are associated with POHC and PO₃-EC defects [27], respectively. The PO₃-EC's signals decrease when CeO₂ was doped into host glasses, which is due to the decreasing precursors of PO₃-EC defects that associated with the NBO in glasses. With the doping of Si in Ce doped sample, the corresponding signals become apparent. These results agree fairly well with the results in Figs. 5 and 6. Besides, these signals of POHC defects were significantly enhanced when these samples were exposed to gamma radiation, as shown the red line in Fig. 7. This is due to the increase of the concentration of POHC defects caused by gamma radiation. Besides, it can be found that the signal is weaker in co-doping samples, especially for the Ce and Sb co-doped glasses. These results indicate that the doping of Si and Sb can decrease the concentration of POHC defects in the phosphate glasses, and the Ce and Sb co-doped sample (G: Ce, Sb) is more apparent.

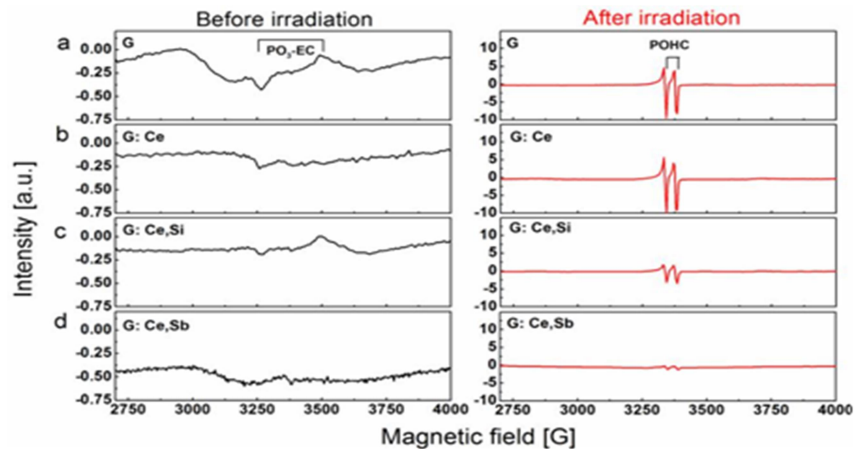


Fig. 7. Measured EPR spectra of a (G), b (G: Ce), c (G: Ce,Si) and d (G: Ce,Sb) glasses before and after gamma radiation (250 krad (Si)).

Most of cerium ions in these optical glasses exists as Ce^{4+} that can be converted to Ce^{3+} by capturing the radiation-induced electrons, and decrease the concentration of $\text{PO}_3\text{-EC}$ defects. Besides, the doping of Ce causes the decrease of NBO, resulting in the decrease of the precursors of $\text{PO}_3\text{-EC}$ and POHC defects. Meanwhile, some Ce^{3+} ions are converted to Ce^{4+} by capturing the radiation-induced holes and then inhibit the formation of POHC defects. In other words, Ce^{3+} and Ce^{4+} ions act as buffer agent to absorb energy, and then give it out by releasing and capturing the electrons and holes to prevent from damaging in the process of gamma radiation. Although the co-doping with Sb_2O_3 in Ce single-doped glass induced the increase of the precursors of POHC defects, Sb^{3+} can be easily photo-oxidized to Sb^{4+} during the gamma radiation by capturing the holes in glasses [28], and Sb^{4+} can be converted to Sb^{5+} by capturing holes in order to improve their stability. The significant improvement of gamma radiation resistance in CeO_2 doped phosphate glass by co-doping with Sb_2O_3 has been accomplished. The radiation resistance, in 385 nm, 525 nm and the broad spectrum (>450 nm), increase dramatically, and it, meanwhile, were distinctly restrained at 1064 nm and 1550 nm, which shows potential applications of these anti-radiation phosphate glasses in the fields of space-born star camera systems, laser window optics, fiber gyroscopes and communications, etc.

4. Conclusion

To conclude, we have shown that the doping of Sb_2O_3 increased the concentration (which is close to the phosphate host glass) of NBO in glass, and weakened the influence of glass structural changes caused by single doping of Ce. Meanwhile, the addition of Sb_2O_3 to Ce-doped optical glasses contributed to a strong improvement of radiation resistance in the visible range, especially at the investigated 525 nm and 385 nm. With further increase of the Ce doping concentration, the radiation resistance was significantly enhanced. In addition, the induced optical loss was also weakened at long wave region, especially for 1064 nm and 1550 nm, which has significant potential applications in space exploring and communications.

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